state is 0 and the energy of the excited state is E_0 .

What is the partition function for this system? Note $\beta = \frac{1}{4}$.

A. $B. Z = 1 + βE_0$ $C. Z = 1 - \beta E_0$ D. None of the above E. ??? $Z = 1 + e^{-\beta E_0}$

- Suppose we have a two-state system, where the energy of the ground
	- 1 *kT*

state is 0 and the energy of the excited state is E_0 .

A. *kT* = 0 B. $kT =$ (some value between 0 and E_0) $c.$ $kT = E_0$

 $D.$ $kT = +\infty$

the ground state?

- Suppose we have a two-state system, where the energy of the ground
- At what temperature is there a 50% probability that the system is in

A system has two states of energies $-\epsilon$ and 2ϵ . What is the probability of observing it in the higher energy state at temperature T?

A. 0
\nB.
$$
\left[1 + \exp\left(\frac{3\epsilon}{kT}\right)\right]^{-1}
$$

\nC. $\left[1 - \exp\left(\frac{3\epsilon}{kT}\right)\right]^{-1}$
\nD. $\left[\exp\left(\frac{\epsilon}{kT}\right) + \exp\left(\frac{-2\epsilon}{kT}\right)\right]^{-1}$
\nE. 1

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What is the partition function of a one-dimensional quantum harmonic oscillator?

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state is 0 and the energy of the excited state is E_0 .

- B. (1 − 1/*e*)*E*⁰
- C. *E*0
- D. None of the above
- E. ???

- Suppose we have a two-state system, where the energy of the ground
- What is the expected value of energy for this system at $kT = E_0$?

A. (1/*e*)*E*⁰

- A. What is the expected value of energy for this particle? B. What is the total energy of the system of N dipoles?
-
- C. Redo part (A) with the partition function $Z_N = Z_1^N$, and show that you get the result that you got in part (B). $Z_N = Z_1^N$
- D. I lied in part (C). Actually, the partition function is $Z_N = \frac{Z_N}{M!} Z_1^N$ if spins are indisguishable. Show that you get the same result as in 1 *N*! Z_1^N 1

Try It Yourself Given the single-particle partition function $Z_1 = e^{-\beta \mu B} + e^{-\beta \mu B}$ $\bar{E} =$ ∑*i e*−*βEi*

both parts (B) & (C).

$$
\frac{\sum_{i} E_{i} e^{-\beta E_{i}}}{\sum_{i} e^{-\beta E_{i}}} = \frac{1}{Z} \sum_{i} E_{i} e^{-\beta E_{i}} = \frac{1}{Z} \left(-\frac{\partial Z}{\partial \beta} \right) = -\frac{\partial \ln Z}{\partial \beta}
$$

$$
e^{+\beta\mu B} = 2\cosh(\beta\mu B)
$$

 $\bar{E} =$

in rotational mode(s) for an CO molecule at room temperature? $\epsilon \approx 2 \times 10^{-4}$ eV

(last slide): *Z*rot ≈ *kT ϵ* $\sum_i E_i e^{-\beta E_i}$ $\sum_i e^{-\beta E_i}$ = 1 *Z* ∑ *i* $E_i e^{-\beta E_i} =$ $\frac{1}{Z}\left(-\frac{\partial Z}{\partial \beta}\right) = -\frac{\partial \ln Z}{\partial \beta}$

A. (1/2)*kT* B. *kT* C. 2*kT*D. None of the above E. ???

For a CO molecule, $\epsilon \approx 2 \times 10^{-4}$ eV. What is the expected energy

- A. Quantum mechanical systems have discrete energy levels, which are not quadratic.
- B. Quantum mechanical systems have relatively big energy-level spacings (compared to kT), which make it a bad approximation to replace the sum with an integral in the proof.
- C. The proof only applies to particles with can be considered disguishable. At low enough temperatures where we can't ignore QM, the fact that we have bosons vs. fermions becomes important.

D. ???

The preceding proof of the equipartition theorem does not apply to

quantum-mechanical systems. Why not?

- A. 0
- B. The most probable speed, moving in the $\pm z$ direction. ±*z*
- C. The rms speed, moving in the $\pm z$ direction. ±*z*
- D. The average speed, moving in the $\pm z$ direction. ±*z*
- E. ???

Suppose you are able to look down the z-axis and view particles following, which is the most likely speed that you would measure?

that are either moving towards you or away from you. Of the

A. B. C. D. E. $(\nu) \propto e^{-\beta m v^2/2}$ (*v*) ∝ *e*−*βmv* $(\nu) \propto \nu e^{-\beta m v^2/2}$ (*v*) ∝ *ve*−*βmv* $(\nu) \propto \nu^2 e^{-\beta m \nu}$

Suppose we lived in 2 spatial dimensions. What would be the form

of the distribution function (probability per unit speed inverval)?

Match the following with the required computation (the limits of the integrals are not specified since they depend on the particular quantity)

A. [∫] *^v* (*v*) *dv* P . $v\mathscr{D}(v) dv$ C. D. $v^2\mathcal{D}(v)dv$ *d dv* (*v*)

- 1. *v*rms
- 2. *v*most probable
- 3. *v*avg
- 4. *v*median

particles). Is this Z equal to Z_1 or Z_N ?

A. Z1 B. ZN C. ???

The previous slide wrote that $S = \frac{1}{T} + k \ln Z$ (for a system of N *U T* + *k* ln *Z*

E. None of the above (or not sure)

From the thermodynamic identity for F (dF $=$...), find an expression

A. *P* = − *kT* $P = +kT$ ∂ ln *Z* ∂*V* ∂ ln *Z* ∂*V*

 $C.$ $P = -kT$ P = $+kT$ ∂*Z* ∂*V* ∂*Z* ∂*V*

for the pressure of a system in terms of the partition function.

of-magnitude of the quantum volume V_Q ?

- A. 10-15 m³.
- B. 10-23 m3 .
- C. 10-31 m3 .
- D. 10-38 m³.
- E. ???

For the nitrogen in this room ($V \sim 100$ m³), what is the correct order-

- A. There is no degeneracy in a 3D particle-in-a-box.
- B. The degeneracy is in the energy, but since we're summing over triplets of n-values and not energy levels, there's no issue.
- C. The degeneracy is a small factor that won't matter for the partition function.
- D. ???

There was no degeneracy factor in the previous expression. Why is

this?

Which of the following will help in the solution to Problem 6.47?

Problem 6.47. Estimate the temperature at which the translational motion of a nitrogen molecule would freeze out, in a box of width 1 cm.

A. This is the temperature at which $V=V_{\mathcal{Q}}.$

- B. This is the temperature where kT is equal to the groundstate energy of a particle-in-a-box.
- C. Both of the above are valid and will get the same answer (to the correct order-of-magnitude).
- D. Neither of the above is the right approach for this problem.

